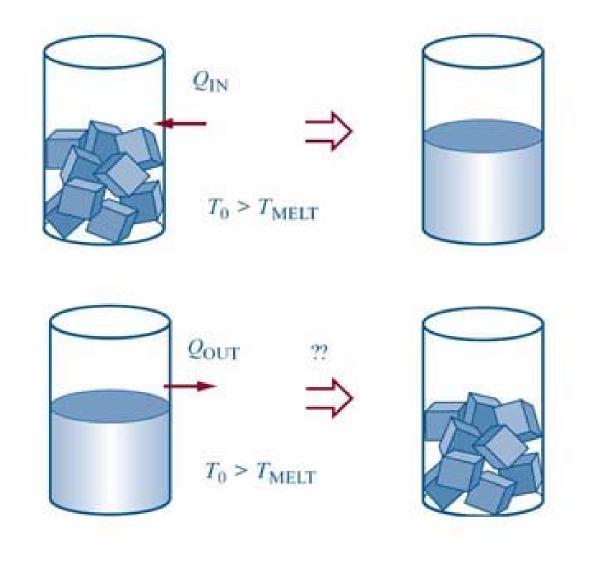
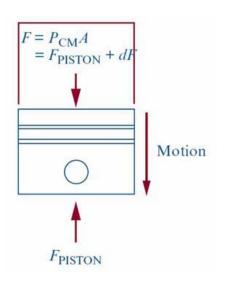
The Second Law of Thermodynamics

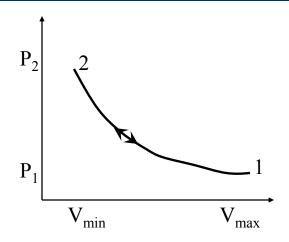
A reversible process

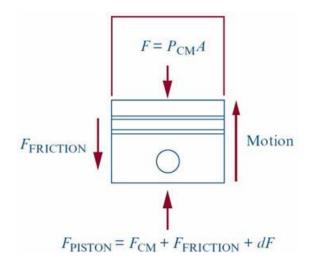


A reversible process ...

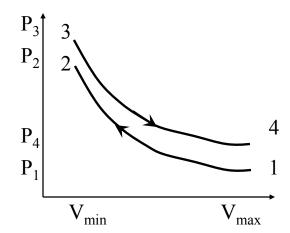


reversible





irreversible



A reversible process ...

- proceeds slowly through equilibrium states.
- could be reversed with no change in heat or work output.

All real processes ...

- are irreversible and have additional heat losses (e.g. due to friction).
- have efficiency $\varepsilon < \varepsilon_{ideal}$

Second Law of Thermodynamics

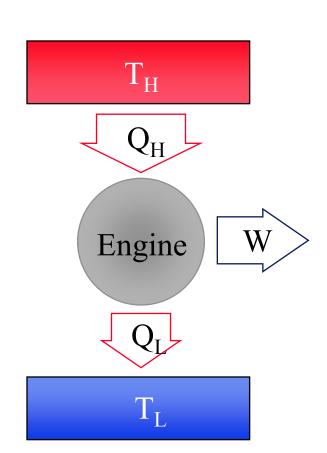
Heat flows spontaneously from a hot object to a cold object, but will not flow spontaneously from a cold object to a hot object.

It is relatively easy to produce thermal energy by doing work (e.g. against friction).

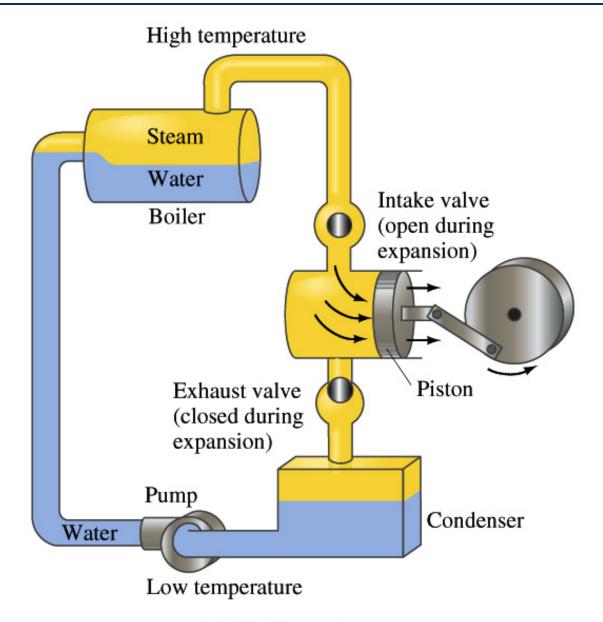
It is also possible to convert internal energy to work.

Heat Engines

- A <u>heat engine</u> converts heat into work.
 - T_H = temperature of heat source
 - T_L = temperature of heat sink
 - Q_H = heat supplied
 - Q_I = heat released
 - W = work produced



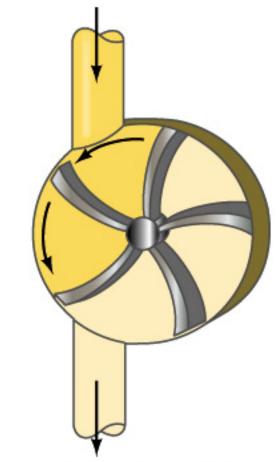
Heat Engines (Steam Engines)



(a) Reciprocating type

Heat Engines (Steam Engines)

High-pressure steam, from boiler



Low-pressure steam, exhausted to condenser

Efficiency of Heat Engines

Conservation of Energy (1st Law):

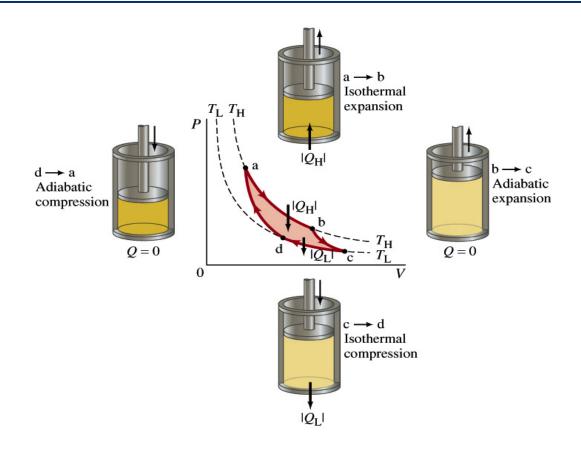
$$0 = \sum Q - \sum W$$
 or $W = Q_H - Q_L$

An engine operates in a cycle, then efficiency is given by

$$\varepsilon = \frac{W}{Q_{H}} = \frac{Q_{H} - Q_{L}}{Q_{H}} = 1 - \frac{Q_{L}}{Q_{H}}$$

 $Q_L > 0$ implies $\varepsilon < 1$

Carnot (Ideal) Heat Engine



Operates in a *reversible* cycle:

a—b: isothermal expansion $(\Delta T = 0)$

b–c: adiabatic expansion (Q=0)

c–d: isothermal compression ($\Delta T=0$)

d–a: adiabatic compression (Q=0)

Ideal (Carnot) Efficiency

$$\frac{Q_L}{Q_H} = \frac{T_L}{T_H} \to \varepsilon_C = 1 - \frac{T_L}{T_H}$$

Efficiencies of Real Heat Engines

- No heat engine can ever have an efficiency greater than that of the Carnot (ideal) heat engine.
- All real heat engines have losses (e.g. friction) and are therefore **not reversible**.
- All real heat engines have efficiencies less than that of a Carnot engine operating between the same temperatures T_H and T_L .

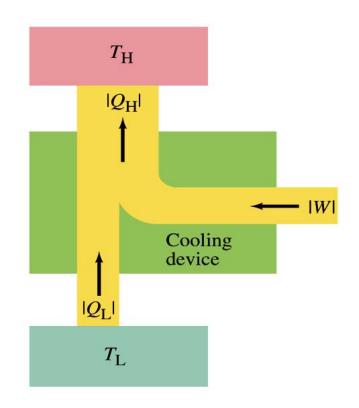
Alternate Statement of Second Law:

• No device can transform a quantity of heat completely to work.

Heat Pump

- A heat pump is a heat engine operating in reverse.
- Examples of heat pumps are refrigerators and air conditioners.
- Conservation of Energy:

Energy In = Energy Out
$$Q_L + W = Q_H$$



$$CP = \frac{Q_L}{W} = \frac{Q_L}{Q_H - Q_L}$$

$$CP_{ideal} = \frac{T_L}{T_H - T_L}$$

Entropy

- Processes that do not violate the first law of thermodynamics (conservation of energy) will never occur spontaneously.
- Entropy (S) is a measure of the disorder or randomness in a system, and is a state variable (like P, V, T) that does not depend on the path taken.

$$dS = \frac{dQ}{T}$$
 (change in entropy S)

where dQ is an infinitesimal heat flow

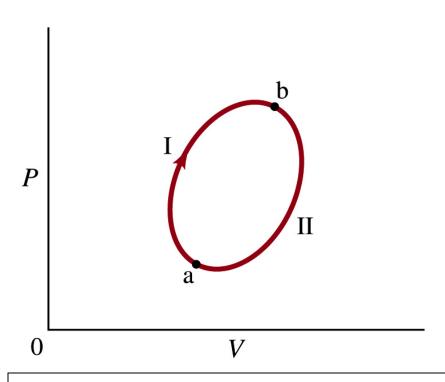
Entropy and Reversible Processes

For any **reversible process** (e.g. Carnot cycle):

$$\oint \frac{dQ}{T} = 0$$

$$\int_{Ia}^{b} \frac{dQ}{T} + \int_{IIb}^{a} \frac{dQ}{T} = 0$$

$$\int_{Ia}^{b} \frac{dQ}{T} = -\int_{IIb}^{a} \frac{dQ}{T} = \int_{IIa}^{b} \frac{dQ}{T}$$



The entropy of a system in a given state is independent of the path taken to get there, and is thus a **state variable**.

 The entropy difference between two equilibrium states a and b does not depend on how the system got from a to b.

$$\Delta S = S_b - S_a = \int_a^b dS = \int_a^b \frac{dQ}{T}$$

Entropy is a **state variable** (like P, V and T)

Second Law in Terms of Entropy

 $\Delta S = 0$ reversible process

 $\Delta S > 0$ irreversible process

Example: Calorimetry $Q = m c \Delta T$

For small changes: dQ = m c dT

$$dS = \frac{dQ}{T}$$

$$\Delta S = \int_{T_i}^{T_f} \frac{dQ}{T} = mc \int_{T_i}^{T_f} \frac{dT}{T} = mc \ln\left(\frac{T_f}{T_i}\right)$$